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## Electronic Effects on Atom Tunneling: Conformational Isomerization of Monomeric *Para*-Substituted Benzoic Acid Derivatives

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Abstract: We present the first generation and spectroscopic identification of the higher-lying E conformer of the simplest aromatic carboxylic acid, benzoic acid (1a), as its O-deuterated isotopologue (E)- $d_1$ -1a using matrix-isolation techniques; the parent (E)-1a could not be observed because of fast H-tunneling to the more stable conformer (Z)-1a. Even deuterated (E)- $d_1$ -1a converts quickly back to (Z)- $d_1$ -1a through D-tunneling with a halflife ( $\tau$ ) of ~12 min in Ar at 11 K. Tunneling computations using an Eckart barrier in conjunction with a CCSD(T)/cc-pVTZ//MP2/ cc-pVDZ + ZPVE intrinsic reaction path revealed that  $\tau$  of (E)-1a is only  $\sim 10^{-5}$  min, in marked contrast to those of simple aliphatic acids, which are in the range of minutes. The electronic substituent effects on D-tunneling in para-substituted benzoic acid derivatives (p-X-PhCOOD,  $d_1$ -1) were systematically studied in Ar matrices at 11 K to derive the first Hammett relationships for atom tunneling.  $\sigma$ -Electron donors (X = alkyl) increase the halflife of  $d_1$ -1, while  $\sigma$ -acceptor/ $\pi$ -donor groups (X = OD, NH<sub>2</sub>, halogen) and to an even greater extent a  $\sigma$ -/ $\pi$ -acceptor group (X = NO<sub>2</sub>) decrease  $\tau$ . The latter finding is in line with the smaller E-to-Z reaction barriers and narrower reaction widths for the isomerization. Tunneling substituent constants ( $\sigma^{t}$ ) for this conformational isomerization were derived experimentally and computationally.

Although double hydrogen transfer in benzoic acid (1a) dimers is a fundamental model for studying the molecular dynamics arising from quantum tunneling in complex biological systems,<sup>1,2</sup> there has been no report regarding the conformational isomerizations through tunneling in *monomeric* 1a; only few simple carboxylic acids have been investigated.<sup>2–4</sup> The lack of systematic studies of electronic substituent effects on tunneling processes motivated us to study the E/Z isomerizations and the potential H- and D-tunneling in 1 (Scheme 1), especially because the electronic effects in *para*-substituted benzoic acids are one of the pillars upon which our understanding of organic reactivity is built (e.g., through the Hammett equation<sup>5</sup>).





The E/Z conformers of carboxylic acids interconvert through C–O bond rotations (Scheme 1), with the Z conformer generally being preferred over a very large temperature range; very few structures, such as glyoxylic acid<sup>6a</sup> and pyruvic<sup>6b</sup> acid, prefer the *E* conformation. While (*E*)-HCOOH [(*E*)-**2**] was first characterized

utilizing microwave spectroscopy in the 1970s,<sup>7a</sup> the first IR signatures for (*E*)-HCOOH<sup>7b</sup> and (*E*)-CH<sub>3</sub>COOH<sup>8</sup> [(*E*)-**3**] trapped in solid matrices were published only in 1997 and 2003, respectively. In contrast to some of the simple aliphatic acids, the *E* isomers of the aromatic acids have not been studied, although several microwave and IR studies of monomeric (*Z*)-benzoic acid [(*Z*)-**1a**] are available.<sup>9</sup> Noticeably, both trapped (*E*)-**2**<sup>4,10</sup> and (*E*)-**3**<sup>11,12</sup> interconvert to their more stable *Z* isomers through H-tunneling. The lifetime of (*E*)-**2** in solid Ar at 8 K is on the order of minutes; under the same conditions, (*E*)-**3** decays ~10 times faster.<sup>8,12</sup>

The characteristic IR bands for (*Z*)-**1a**, especially the single O–H stretching vibration (3570.5 cm<sup>-1</sup>, Ar, 11 K) confirmed the presence of only one conformer. We attempted to prepare and identify the higher-lying (*E*)-**1a** either photochemically by irradiation of (*Z*)-**1a** at  $\lambda = 254$  nm (**1a** has absorption maxima at  $\lambda = 200$ , 230, and 280 nm<sup>13</sup>) for 30 min in matrices at 11 K or thermally by flowing (*Z*)-**1a** through an 800 °C quartz glass tube followed by trapping in various matrices at 11 K. Although our CCSD(T)/cc-pVTZ// MP2/cc-pVDZ computations gave an isomerization barrier of 5.0 kcal mol<sup>-1</sup> (structures optimized using density functional theory at various levels gave similar energy differences; see the Supporting Information), which cannot be overcome thermally at 11 K, we were unable to detect (*E*)-**1a**. The same applies to the *para*-substituted derivatives **1b**–**f**, in stark contrast to the previously observed matrix-isolated (*E*)-**2** and (*E*)-**3**.

To examine whether the isomerization through H-tunneling from (*E*)-**1a** to (*Z*)-**1a** is simply too fast on the time scale of our experiments (minutes) or we had not generated the (*E*)-**1a** isomer, we examined  $d_1$ -**1a** (OD deuteration) in the same way. Indeed, we were able to record the IR spectrum of (*E*)- $d_1$ -**1a** in Ar at 11 K (Figure 1); its fundamental frequencies  $v_{O-D}$ ,  $v_{C=O}$ , and  $\delta_{C-H,in-plane}$  were blue-shifted and the  $v_{C-O}$  and  $\delta_{C-H,out-of-plane}$  absorptions red-shifted relative to those of (*Z*)- $d_1$ -**1a**.



**Figure 1.** Difference IR spectra of (E)- $d_1$ -**1a** (upward-pointing peaks) and (Z)- $d_1$ -**1a** (downward-pointing peaks): (a) computed at the MP2/cc-pVDZ level (unscaled); (b) measured in solid Ar at 11 K after 30 min of irradiation at 254 nm. The splitting of the  $\nu_{C-O}$  band of (Z)- $d_1$ -**1a** is due to matrix effects.

Surprisingly,  $(E)-d_1-1a$  did not persist under these conditions: it converted into  $(Z)-d_1-1a$  with a half-life  $(\tau)$  of 12 min in Ar (Table 1). The relative temperature independence of the half-lives at 11 and 20 K and the apparently very large primary H/D kinetic isotope effect support the notion of a tunneling mechanism. The D-tunneling of the  $(E)-d_1-1a$  isomer is three orders of magnitude faster than that of  $(E)-d_1-2$ , with a half-life of  $\sim 7$  days<sup>4</sup> in Ar at 4.3 K. Clearly, the phenyl ring has a decisive electronic effect on the tunneling rate, prompting us to study a selection of *p*ara-substituted, monodeuterated benzoic acids  $(d_1-1b-f)$ .

The matrix material influences the tunneling rates significantly (Table 1), a trend also seen in other tunneling processes.<sup>4,14</sup> The rate retardation in N<sub>2</sub> matrices is likely due to formation of a hydrogen-bond complex between the acid function and N<sub>2</sub>. Still, the stabilization of the (*E*)-**1a** rotamer in N<sub>2</sub> was not sufficient to allow its spectroscopic observation.

**Table 1.** Half-Lives  $\tau$  (min)<sup>a</sup> of (*E*)- $d_1$ -**1a** in Various Matrices

<i>T</i> (K)	Ar	Xe	N <sub>2</sub>
11 20	$\begin{array}{c} 12\pm2\\ 12\pm2 \end{array}$	$\begin{array}{c} 22\pm3\\ 25\pm3 \end{array}$	$11 \pm 1$ daysb

<sup>*a*</sup> From first-order kinetics measurements based on the slope  $[k (s^{-1})]$  of a plot of ln(C=O peak height) vs time. <sup>*b*</sup> Not determined because of decomposition of the N<sub>2</sub> matrix at this temperature.

The tunneling process was also analyzed through computation of the intrinsic reaction path (IRP) connecting the rotational transition structure (TS-1a) with (E)-1a and (Z)-1a. A final potential energy curve along the isomerization IRP was then constructed from CCSD(T)/cc-pVTZ energy points and MP2/cc-pVDZ zero-point vibrational energies (ZPVEs). Tunneling probabilities were evaluated using an unscaled asymmetric Eckart potential<sup>15</sup> to the imaginary TS-1a barrier frequency  $\omega^* = 536i \text{ cm}^{-1}$  (394*i* cm<sup>-1</sup> for TS- $d_1$ -1a) as well as the ZPVE-corrected reaction energy of -6.0 kcal mol<sup>-1</sup> (-6.0 kcal mol<sup>-1</sup> for  $d_1$ -1a) and barrier height of +5.0 (+5.3 kcal mol<sup>-1</sup> for  $d_1$ -1a). This procedure has been shown to yield tunneling half-lives in good qualitative agreement with experiment for the isotopologues of hydroxycarbene<sup>14</sup> and phenylhydroxycarbene.<sup>16</sup> A vibrational "reaction" mode of (E)-1a toward TS-1a with a frequency of  $\omega_0 = 498 \text{ cm}^{-1}$  [358 cm<sup>-1</sup> for (E)- $d_1$ -1a] was identified. In good qualitative agreement with the experiment, the computations gave a half-life of 2.8 h (55 min at the MP2/cc-pVDZ level) for (*E*)- $d_1$ -**1a**, while they predicted  $\tau \sim$  $10^{-5}$  min for the H-tunneling in (E)-1a, which is much too fast to be measured using our present experimental setup. The noted much faster D-tunneling in  $d_1$ -1a relative to  $d_1$ -2<sup>4</sup> can be rationalized by the lower reaction barrier [by 2.4 kcal mol<sup>-1</sup> at the CCSD(T)/ccpVTZ//MP2/cc-pVDZ level] and especially the much smaller barrier width (Figure 2).

The experimental and computed half-lives for *para*-substituted benzoic acids (Table 2) revealed the preliminary trend that pure  $\sigma$ -electron-donating groups (CH<sub>3</sub>, *t*-Bu) increase the tunneling half-life of (*E*)-*d*<sub>1</sub>-**1**, while  $\sigma$ -acceptor/ $\pi$ -donor groups (OD, NH<sub>2</sub>, halogen) and to an even greater extent a  $\sigma$ -/ $\pi$ -electron-acceptor group (NO<sub>2</sub>) decrease the D-tunneling half-life; the *d*<sub>1</sub>-**1f** analogue had the lowest and narrowest IRP curve among those for *d*<sub>1</sub>-**1** (Figure 2). However, a Hammett plot showed that the substituent constant  $\sigma$  does not correlate well with  $\tau$  because the tunneling isomerization in **1** apparently is electronically too different from that for simple *para*-substituted benzoic acid ionization (Figure S27). Hence, we specifically derived the tunneling substituent constants  $\sigma^{t}$  (Table 2) using the relation  $\rho\sigma^{t} = pk_{\rm H} - pk_{\rm X} = \log(k_{\rm X}/k_{\rm H})$ ,<sup>5</sup> with the reaction constant  $\rho$  set equal to 1; the excellent fit is



**Figure 2.** E/Z isomerization IRPs for **1a**,  $d_1$ -**1a**-**f**, **2**, and  $d_1$ -**2** computed at the MP2/cc-pVDZ level.

**Table 2.** Experimental (Ar, 11 K) and Computed (MP2/cc-pVDZ, 11 K; in parentheses) Half-Lives  $\tau$  of (*E*)-*d*<sub>1</sub>-**1a**-**f** and Derived  $\sigma$ <sup>t</sup> Values for the Tunneling *E*/*Z* Isomerizations

	Х	au (min)	$\sigma^{t}$
( <i>E</i> )- <i>d</i> <sub>1</sub> - <b>1b</b>	CH <sub>3</sub>	$25 \pm 2$ (188)	-0.38 (-0.53)
$(E)-d_1-1c$	t-Bu	$16 \pm 2 (140)$	-0.17(-0.40)
(E)-d <sub>1</sub> -1a	H	$12 \pm 2 (55)$	0.00 (0.00)
$(E)-d_1-1d$	OD	$6 \pm 2 (28^{a})$	+0.26(+0.30)
$(E)-d_1-1e$	Cl	$3 \pm 1$ (8)	+0.60(+0.85)
$(E)$ - $d_1$ - <b>1f</b>	$NO_2$	$3 \pm 1$ (3)	+0.53 (+1.26)

<sup>&</sup>lt;sup>*a*</sup> The two C–OD*para* conformers showed slightly different half-lives (for details, see the Supporting Information).

depicted in Figure S28 in the Supporting Information. Importantly, the signs of  $\sigma$  and  $\sigma^{t}$  were the same, even when the absolute values were quite different. Figure 3 underlines these qualitative findings by correlating the computed barrier widths with the experimental half-lives.



*Figure 3.* Plot of the barrier widths derived from IRP computations at the MP2/cc-pVDZ vs the experimental half-live of (E)- $d_1$ -1.

A natural extension of our studies is the examination of conformational tunneling in biologically relevant systems such as amino acids and fatty acids. Comprehensive computational and experimental studies on a generalization of the electronic substituent effects on tunneling processes are in progress.

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**Supporting Information Available:** Detailed experimental and computational data. This material is available free of charge via the Internet at http://pubs.acs.org.

## References

- (1) (a) Horsewill, A. J. J. Phys. Org. Chem. 2010, 23, 580–585. (b) Horsewill, A. J.; McGloin, C. J.; Trommsdorff, H. P.; Johnson, M. R. Chem. Phys. 2003, 291, 41–52. (c) Rambaud, C.; Oppenländer, A.; Pierre, M.; Trommsdorff, H. P.; Vial, J.-C. Chem. Phys. 1989, 136, 335–347. (d) Watt, C. I. F. J. Phys. Org. Chem. 2010, 23, 561–571.
- (2) Fillaux, F.; Romain, F.; Limage, M.-H.; Leygue, N. Phys. Chem. Chem. Phys. 2006, 8, 4327–4336.
- (3) (a) Trakhtenberg, L. I.; Fokeyev, A. A.; Zyubin, A. S.; Mebel, A. M.; Lin,
   S. H. J. Chem. Phys. 2009, 130, 144502/1–11. (b) Schomaker, V.;
   Ogorman, J. M. J. Am. Chem. Soc. 1947, 69, 2638–2644. (c) Karle, I. L.;
   Brockway, L. O. J. Am. Chem. Soc. 1944, 66, 1974–1979.
- (4) Domanskaya, A.; Marushkevich, K.; Khriachtchev, L.; Räsänen, M. J. Chem. Phys. 2009, 130, 154509/1-5.
- (5) (a) Hammett, L. P. J. Am. Chem. Soc. 1937, 59, 96–103. (b) Jones, R. A. Y. Physical and Mechanistic Organic Chemistry: Cambridge University Press: Cambridge, U.K., 1979. (c) Anslyn, E. V.; Dougherty, D. A. Modern Physical Organic Chemistry: University Science Books: Sausalito, CA, 2006.
- Organic Chemistry; University Science Books: Sausalito, CA, 2006.
  (6) (a) van Eijck, B. P.; van Duineveldt, F. B. J. Mol. Struct. 1977, 39, 157–163. (b) Dyllick-Brenzinger, C. E.; Bauder, A.; Günthard, H. H. Chem. Phys. 1977, 23, 195–206.
- (7) (a) Hocking, W. M. Z. Naturforsch. A 1976, 31, 1113–1121; Bjarnov, E.; Hocking, W. M. Z. Naturforsch. A 1978, 33, 610–618. (b) Pettersson, M.; Lundell, J.; Khriachtchev, L.; Räsänen, M. J. Am. Chem. Soc. 1997, 119, 11715–11716.

- (8) Macoas, E. M. S.; Khriachtchev, L.; Pettersson, M.; Fausto, R.; Räsänen, M. J. Am. Chem. Soc. 2003, 125, 16188–16189.
- (9) (a) Onda, M.; Asai, M.; Takise, K.; Kuwae, K.; Hayami, K.; Kuroe, A.; Mori, M.; Miyazaki, H.; Suzuki, N.; Yamaguchi, I. J. Mol. Struct. 1999, 482, 301–303. (b) Reva, I. D.; Stepanian, S. G. J. Mol. Struct. 1995, 349, 337–340.
- (10) (a) Pettersson, M.; Macoas, E. M. S.; Khriachtchev, L.; Fausto, R.; Räsänen, M. J. Am. Chem. Soc. 2003, 125, 4058–4059. (b) Pettersson, M.; Macoas, E. M. S.; Khriachtchev, L.; Lundell, J.; Fausto, R.; Räsänen, M. J. Chem. Phys. 2002, 117, 9095–9098.
- (11) (a) Macoas, E. M. S.; Khriachtchev, L.; Pettersson, M.; Fausto, R.; Räsänen, M. J. Chem. Phys. 2004, 121, 1331–1338. (b) Meyer, R.; Ha, T. K.; Frei, H.; Gunthard, H. H. Chem. Phys. 1975, 9, 393–402. (c) Redington, R. L.; Lin, K. C. J. Chem. Phys. 1971, 54, 4111–4119. (d) Berney, C. V.; Redington, R. L.; Lin, K. C. J. Chem. Phys. 1970, 53, 1713–1721.
- (12) Macoas, E. M. S.; Khriachtchev, L.; Fausto, R.; Räsänen, M. J. Phys. Chem. A 2004, 108, 3380–3389.
- (13) Hosoya, H.; Tanaka, J.; Nagakura, S. J. Mol. Spectrosc. 1962, 8, 257–275.
- (14) Schreiner, P. R.; Reisenauer, H. P.; Pickard, F. C.; Simmonett, A. C.; Allen,
- W. D.; Matyus, E.; Csaszar, A. G. *Nature* **2008**, *453*, 906–909.
- (15) Eckart, C. Phys. Rev. 1930, 35, 1303-1309.
- (16) Gerbig, D.; Reisenauer, H. P.; Wu, C. H.; Ley, D.; Allen, W. D.; Schreiner, P. R. J. Am. Chem. Soc. **2010**, 132, 7273–7275.

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